

teorema di unicità

$$\nabla \times \vec{E}^a = -j\omega\mu \vec{H}^a$$

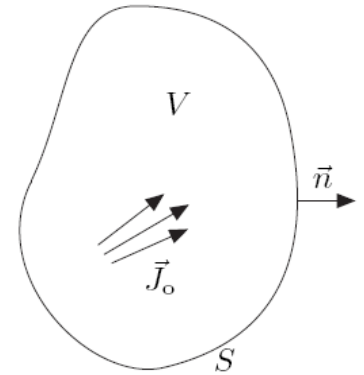
$$\nabla \times \vec{E}^b = -j\omega\mu \vec{H}^b$$

$$\nabla \times \vec{H}^a = j\omega\epsilon \vec{E}^a + \vec{J}_o$$

$$\nabla \times \vec{H}^b = j\omega\epsilon \vec{E}^b + \vec{J}_o$$

$$\delta\vec{E} = \vec{E}^a - \vec{E}^b$$

$$\delta\vec{H} = \vec{H}^a - \vec{H}^b$$



$$\nabla \times \delta\vec{E} = -j\omega\mu \delta\vec{H}$$

$$\nabla \times \delta\vec{H} = j\omega\epsilon \delta\vec{E}$$

$$0 = \underbrace{\frac{1}{2} \omega \int_V (\mu_o \mu'' |\delta\vec{H}|^2 + \epsilon_o \epsilon'' |\delta\vec{E}|^2) dV}_{\delta P_{\text{diss}}} + \underbrace{\frac{1}{2} \text{Re} \left(\int_S \delta\vec{E} \times \delta\vec{H}^* \cdot \vec{n} dS \right)}_{\delta P_{\text{out}}}$$

$$\vec{n} \times \vec{E} - \mathcal{Z} \vec{H}_T = \vec{V}_T \quad \text{in } S_1 \quad \text{Re}(\mathcal{Z}) > 0$$

$$\vec{H} \times \vec{n} - \mathcal{Y} \vec{E}_T = \vec{W}_T \quad \text{in } S_2 \quad \text{Re}(\mathcal{Y}) > 0$$

$$\delta P_{\text{out}} = \frac{1}{2} \text{Re}(\mathcal{Z}) \int_{S_1} |\delta \vec{H}_T|^2 dS + \frac{1}{2} \text{Re}(\mathcal{Y}) \int_{S_2} |\delta \vec{E}_T|^2 dS \geq 0$$

- *l'assegnazione del campo elettrico tangenziale*

$$\vec{n} \times \vec{E} = \vec{V}_T$$

- *l'assegnazione del campo magnetico tangenziale*

$$\vec{H} \times \vec{n} = \vec{W}_T$$

- *la condizione di parete elettrica*

$$\vec{n} \times \vec{E} = 0$$

- *la condizione di parete magnetica*

$$\vec{n} \times \vec{H} = 0$$

- *la condizione d'impedenza*

$$\vec{n} \times \vec{E} = \mathcal{Z} \vec{H}_T$$

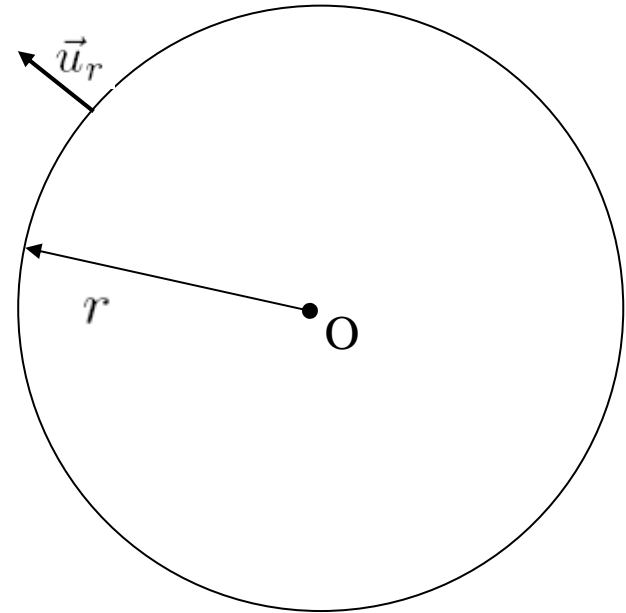
- *la condizione di ammettenza*

$$\vec{H} \times \vec{n} = \mathcal{Y} \vec{E}_T$$

- *la condizione di radiazione*

Condizione di radiazione

$$\left. \begin{aligned} \lim_{r \rightarrow \infty} r^\nu \vec{E} &= 0 \\ \lim_{r \rightarrow \infty} r^\nu \vec{H} &= 0 \end{aligned} \right\} \quad \forall \nu < 1$$
$$\lim_{r \rightarrow \infty} r \left(\vec{u}_r \times \vec{E} - \eta_\infty \vec{H} \right) = 0$$
$$\lim_{r \rightarrow \infty} r \left(\vec{H} \times \vec{u}_r - \vec{E} / \eta_\infty \right) = 0$$



radiazione

$$\vec{A} = \mu \int_V \frac{e^{-jkR}}{4\pi R} \vec{J}_o(\vec{r}') dV'$$

$$\begin{aligned} \vec{E} &= -j\omega \left(\vec{A} + \frac{1}{k^2} \nabla \nabla \cdot \vec{A} \right) \\ &= -j\eta k \int_V \frac{e^{-jkR}}{4\pi R} \left[\left(1 - \frac{j}{kR} - \frac{1}{k^2 R^2} \right) \vec{J}_o(\vec{r}') \right. \\ &\quad \left. - \left(1 - \frac{3j}{kR} - \frac{3}{k^2 R^2} \right) \vec{u}_R \left(\vec{u}_R \cdot \vec{J}_o(\vec{r}') \right) \right] dV' \end{aligned}$$

$$\begin{aligned} \vec{H} &= \frac{\nabla \times \vec{A}}{\mu} \\ &= -jk \int_V \frac{e^{-jkR}}{4\pi R} \left(1 - \frac{j}{kR} \right) \vec{u}_R \times \vec{J}_o(\vec{r}') dV' \end{aligned}$$

$$\vec{u}_R = \frac{\vec{R}}{R}$$

zona di radiazione

$$r \gg D \quad \Rightarrow \quad \frac{1}{R} \approx \frac{1}{r} \quad \vec{u}_R \approx \vec{u}_r$$

$$r > \frac{2D^2}{\lambda} \quad \Rightarrow \quad e^{-jkR} \approx e^{-jkr} e^{j\vec{u}_r \cdot \vec{r}'}$$

$$r \gg \lambda \quad \Rightarrow \quad \frac{1}{k^2 r^2} \ll \frac{1}{kr} \ll 1$$

$$\vec{A} \approx \mu \frac{e^{-jkr}}{4\pi r} \underbrace{\int_V e^{jk\vec{u}_r \cdot \vec{r}'} \vec{J}_o(\vec{r}') dV'}_{\vec{N}(\vartheta, \varphi)}$$

$$\vec{E} \approx -j\eta k \frac{e^{-jkr}}{4\pi r} \left(\vec{N} - \vec{u}_r (\vec{u}_r \cdot \vec{N}) \right)$$

$$\vec{H} \approx \frac{\vec{u}_r \times \vec{E}}{\eta}$$

zona dei campi reattivi

$$r \gg D \quad \Rightarrow \quad \frac{1}{R} \approx \frac{1}{r} \quad \vec{u}_R \approx \vec{u}_r$$

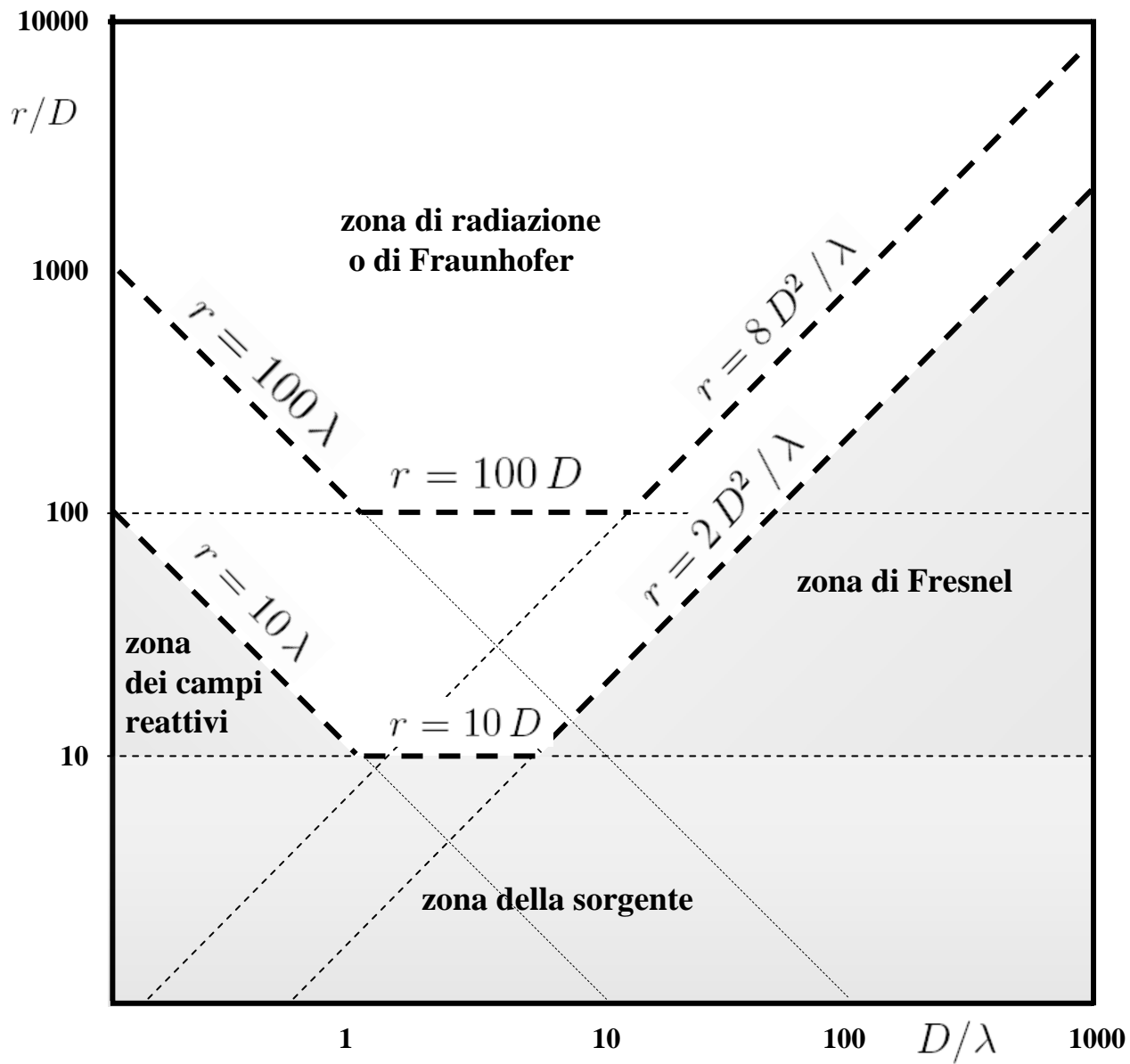
$$r > \frac{2D^2}{\lambda} \quad \Rightarrow \quad e^{-jkR} \approx e^{-jkr} e^{j\vec{u}_r \cdot \vec{r}'}$$

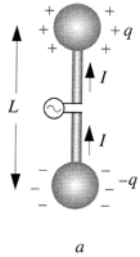
$$r \ll \lambda \quad \Rightarrow \quad 1 \ll \frac{1}{kr} \ll \frac{1}{k^2 r^2}$$

$$\vec{A} \approx \mu \frac{e^{-jkr}}{4\pi r} \underbrace{\int_V e^{jk\vec{u}_r \cdot \vec{r}'} \vec{J}_o(\vec{r}') dV'}_{\vec{N}(\vartheta, \varphi)}$$

$$\vec{E} \approx -j\eta k \frac{e^{-jkr}}{4\pi r} \left[\left(\cancel{1} - \cancel{\frac{j}{kr}} - \frac{1}{k^2 r^2} \right) \vec{N} - \left(\cancel{1} - \cancel{\frac{3j}{kr}} - \frac{3}{k^2 r^2} \right) \vec{u}_r (\vec{u}_r \cdot \vec{N}) \right]$$

$$\vec{H} \approx -jk \frac{e^{-jkr}}{4\pi r} \left(\cancel{1} - \frac{j}{kr} \right) \vec{u}_r \times \vec{N}$$





campo di radiazione del dipolo hertziano

$$(r \gg \lambda)$$

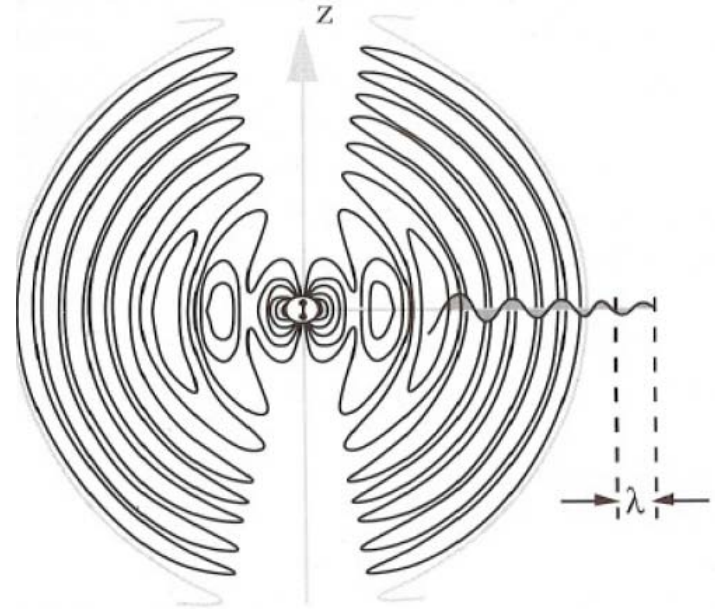
$$\vec{N} = I L \vec{u}_z$$

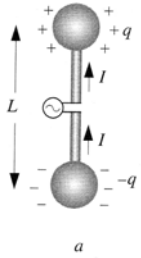
$$\vec{E} = j \eta_0 I L \frac{e^{-jkr}}{2 \lambda_0 r} \sin \vartheta \vec{u}_\vartheta$$

$$\vec{H} = j I L \frac{e^{-jkr}}{2 \lambda_0 r} \sin \vartheta \vec{u}_\varphi$$

$$K = \frac{\eta_0}{8} \left(\frac{L}{\lambda_0} \right)^2 |I|^2 \sin^2 \vartheta$$

$$P_{\text{irr}} = \frac{\eta_0 \pi}{3} \left(\frac{L}{\lambda_0} \right)^2 |I|^2$$





campo reattivo del dipolo hertziano

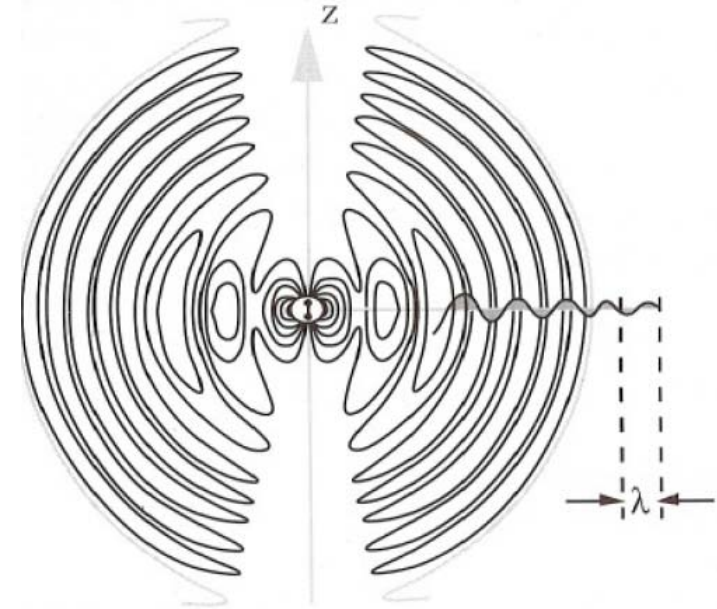
$$(r \ll \lambda, r \gg L)$$

$$\vec{E} \approx \frac{qL}{4\pi\epsilon_0 r^3} (2 \cos \vartheta \vec{u}_r + \sin \vartheta \vec{u}_\vartheta)$$

$$= -\nabla \frac{qL}{4\pi\epsilon_0 r^2} \cos \vartheta$$

$$\vec{H} \approx \frac{IL}{4\pi r^2} \sin \vartheta \vec{u}_\varphi$$

$$= \frac{IL}{4\pi} \frac{\vec{u}_z \times \vec{u}_r}{r^2} \quad (\text{Biot-Savart})$$



$$I = j\omega q$$

$$\frac{\eta_0 \lambda_0}{2\pi} = \frac{1}{\epsilon_0 \omega}$$

spira di piccole dimensioni

$$\vec{N} = j k S I \sin \vartheta \vec{u}_\varphi$$

$$\vec{E} = \eta_o I \frac{\pi S}{\lambda_o^2} \frac{e^{-jkr}}{r} \sin \vartheta \vec{u}_\varphi$$

$$\vec{H} = -I \frac{\pi S}{\lambda_o^2} \frac{e^{-jkr}}{r} \sin \vartheta \vec{u}_\vartheta$$

S = area della spira
(πa^2 spira circolare)

$$K = \frac{\eta_o |I|^2}{2} \left(\frac{\pi S}{\lambda_o^2} \right)^2 \sin^2 \vartheta$$

$$P_{irr} = \frac{\eta_o \pi |I|^2}{3} \left(\frac{2\pi S}{\lambda_o^2} \right)^2$$

dipoli di lunghezza confrontabile con la lunghezza d'onda

$$I(z) = I_o \frac{\sin k(L - |z|)}{\sin(kL)}$$

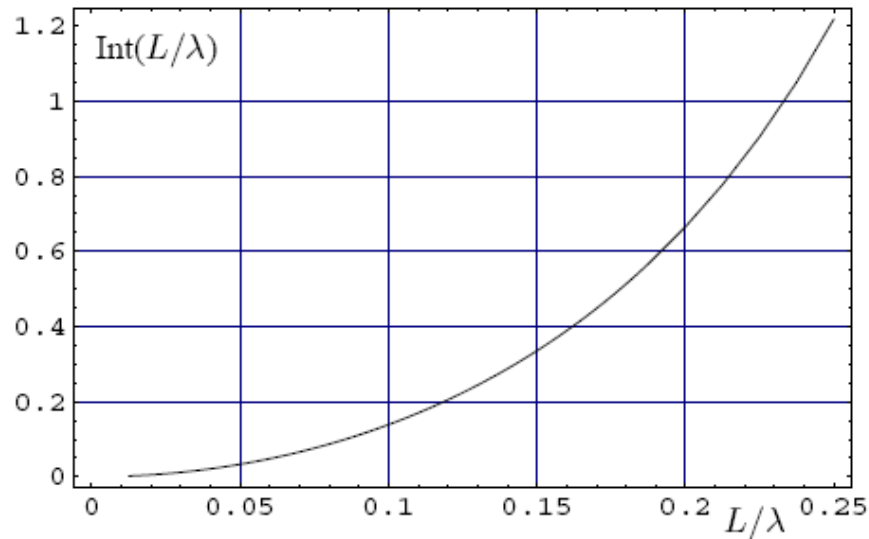
$$\vec{N} = 2 \vec{u}_z I_o \frac{\cos(kL \cos \vartheta) - \cos(kL)}{k \sin(kL) \sin^2 \vartheta}$$

$$\vec{E} = j \eta I_o \frac{e^{-jkr}}{2 \pi r} \frac{\cos(kL \cos \vartheta) - \cos(kL)}{\sin(kL) \sin \vartheta} \vec{u}_\vartheta$$

$$\vec{H} = j I_o \frac{e^{-jkr}}{2 \pi r} \frac{\cos(kL \cos \vartheta) - \cos(kL)}{\sin(kL) \sin \vartheta} \vec{u}_\varphi$$

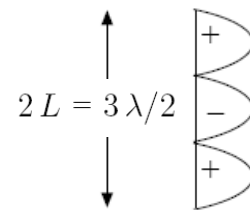
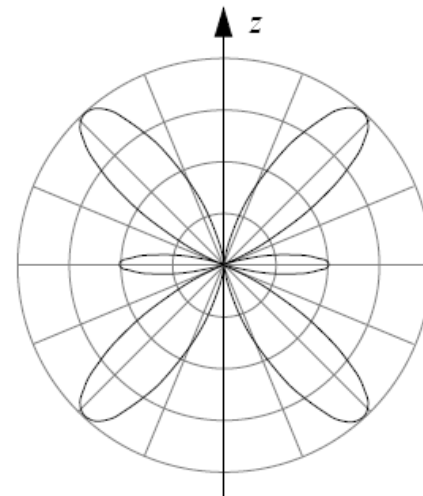
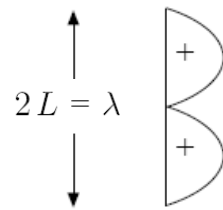
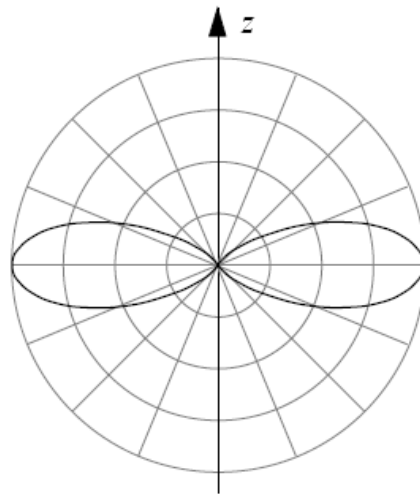
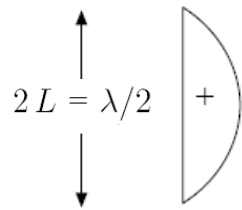
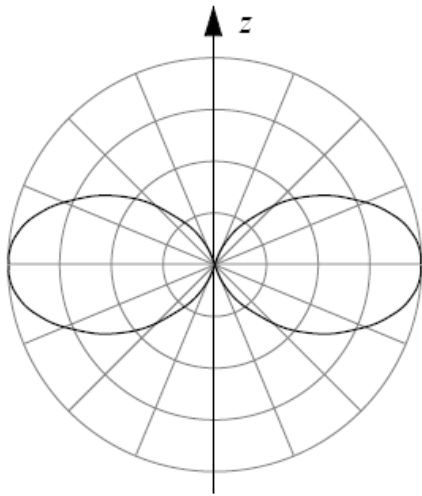
$$K = \frac{\eta}{8 \pi^2} |I_o|^2 \left(\frac{\cos(kL \cos \vartheta) - \cos(kL)}{\sin(kL) \sin \vartheta} \right)^2$$

$$\int_0^{2\pi} \int_0^\pi \left(\frac{\cos(kL \cos \vartheta) - \cos(kL)}{\sin(kL) \sin \vartheta} \right)^2 \sin \vartheta d\vartheta d\varphi = 2\pi \text{Int}(L/\lambda)$$



$$P_{\text{irr}} = \frac{\eta}{4\pi} |I_o|^2 \text{Int}(L/\lambda)$$

$$D = \frac{2}{\text{Int}(L/\lambda)} \left(\frac{\cos(kL \cos \vartheta) - \cos(kL)}{\sin(kL) \sin \vartheta} \right)^2$$



equazioni di Maxwell generalizzate

$$\nabla \times \vec{E} = -j\omega\mu \vec{H} - \vec{M}_o$$

$$\nabla \times \vec{H} = j\omega\epsilon \vec{E} + \vec{J}_o$$

dualità

\vec{E}	\rightarrow	\vec{H}
\vec{H}	\rightarrow	$-\vec{E}$
\vec{J}_o	\rightarrow	\vec{M}_o
\vec{M}_o	\rightarrow	$-\vec{J}_o$
ϵ	\leftrightarrow	μ

sorgenti elettriche

sorgenti magnetiche

$$\nabla \times \vec{E} = -j\omega\mu\vec{H}$$

$$\nabla \times \vec{E} = -j\omega\mu\vec{H} - \vec{M}_o$$

$$\nabla \times \vec{H} = j\omega\epsilon\vec{E} + \vec{J}_o$$

$$\nabla \times \vec{H} = j\omega\epsilon\vec{E}$$

$$\vec{E} = -j\omega \left(\vec{A} + \frac{1}{k^2} \nabla \nabla \cdot \vec{A} \right)$$

$$\vec{H} = -j\omega \left(\vec{F} + \frac{1}{k^2} \nabla \nabla \cdot \vec{F} \right)$$

$$\vec{H} = \frac{1}{\mu} \nabla \times \vec{A}$$

$$\vec{E} = -\frac{1}{\epsilon} \nabla \times \vec{F}$$

$$\vec{A} = \mu \int_V \frac{e^{-jkR}}{4\pi R} \vec{J}_o(\vec{r}') dV'$$

$$\vec{F} = \epsilon \int_V \frac{e^{-jkR}}{4\pi R} \vec{M}_o(\vec{r}') dV'$$

nell'approssimazione di campo di radiazione

sorgenti elettriche

sorgenti magnetiche

$$\vec{A} = \mu \frac{e^{-jkr}}{4\pi r} \vec{N}$$

$$\vec{F} = \epsilon \frac{e^{-jkr}}{4\pi r} \vec{L}$$

$$\vec{N} = \int_V e^{jk\vec{u}_r \cdot \vec{r}'} \vec{J}_o(\vec{r}') dV'$$

$$\vec{L} = \int_V e^{jk\vec{u}_r \cdot \vec{r}'} \vec{M}_o(\vec{r}') dV'$$

$$\vec{E} = -j\eta \frac{e^{-jkr}}{2\lambda r} \left(\vec{N} - \vec{u}_r (\vec{u}_r \cdot \vec{N}) \right)$$

$$\vec{H} = -\frac{j}{\eta} \frac{e^{-jkr}}{2\lambda r} \left(\vec{L} - \vec{u}_r (\vec{u}_r \cdot \vec{L}) \right)$$

$$\vec{H} = -j \frac{e^{-jkr}}{2\lambda r} \vec{u}_r \times \vec{N}$$

$$\vec{E} = j \frac{e^{-jkr}}{2\lambda r} \vec{u}_r \times \vec{L}$$

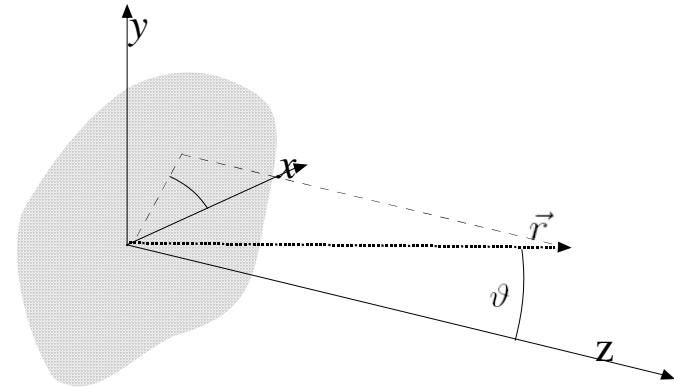
$$K = \frac{\eta}{8\lambda^2} \left(|\vec{N}|^2 - |\vec{u}_r \cdot \vec{N}|^2 \right)$$

$$K = \frac{1}{8\eta\lambda^2} \left(|\vec{L}|^2 - |\vec{u}_r \cdot \vec{L}|^2 \right)$$

$$\vec{H} = \frac{\vec{u}_r \times \vec{E}}{\eta} \quad \vec{E} = \eta \vec{H} \times \vec{u}_r$$

$$K = \frac{|\vec{E}|^2}{2\eta} r^2 = \frac{\eta |\vec{H}|^2}{2} r^2$$

radiazione da aperture



$$\begin{aligned}\vec{L} &= 2 \int_{\text{apertura}} e^{jk \vec{u}_r \cdot \vec{r}'} M_s(\vec{r}') ds' \\ &= 2 \int_{\text{apertura}} e^{jk(u x' + v y')} M_s(x', y') dx' dy'\end{aligned}$$

$$\vec{E} = j \frac{e^{-jkr}}{2 \lambda r} \vec{u}_r \times \vec{L}$$

$$\vec{H} = -\frac{j}{\eta} \frac{e^{-jkr}}{2 \lambda r} \left(\vec{L} - \vec{u}_r (\vec{u}_r \cdot \vec{L}) \right)$$

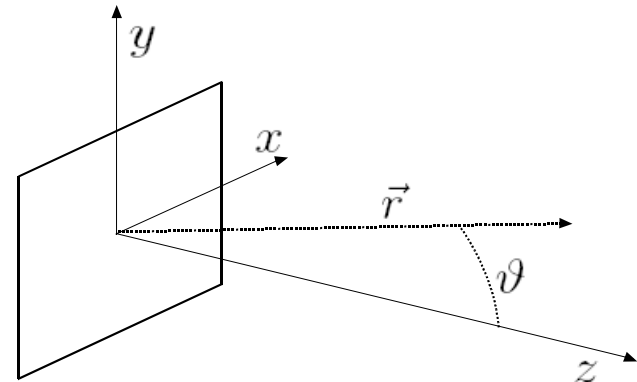
$$K = \frac{1}{8 \eta \lambda^2} \left(|\vec{L}|^2 - |\vec{u}_r \cdot \vec{L}|^2 \right)$$

$$u = \vec{u}_r \cdot \vec{u}_x = \sin \vartheta \cos \varphi$$

$$v = \vec{u}_r \cdot \vec{u}_y = \sin \vartheta \sin \varphi$$

$$\vec{M}_s = \vec{E} \times \vec{u}_z$$

Apertura rettangolare illuminata uniformemente $\vec{E}_o = E_o \vec{u}_x$



$$\begin{aligned}
 \vec{L} &= 2 \int_{\text{apertura}} e^{j k \vec{u}_r \cdot \vec{r}'} M_s(\vec{r}') ds' \\
 &= -2 E_o \vec{u}_y \int_{-a/2}^{a/2} e^{j k u x'} dx' \int_{-b/2}^{b/2} e^{j k v y'} dy' \\
 &= -2 E_o a b \operatorname{sinc}(\pi u a / \lambda) \operatorname{sinc}(\pi v b / \lambda) \vec{u}_y
 \end{aligned}$$

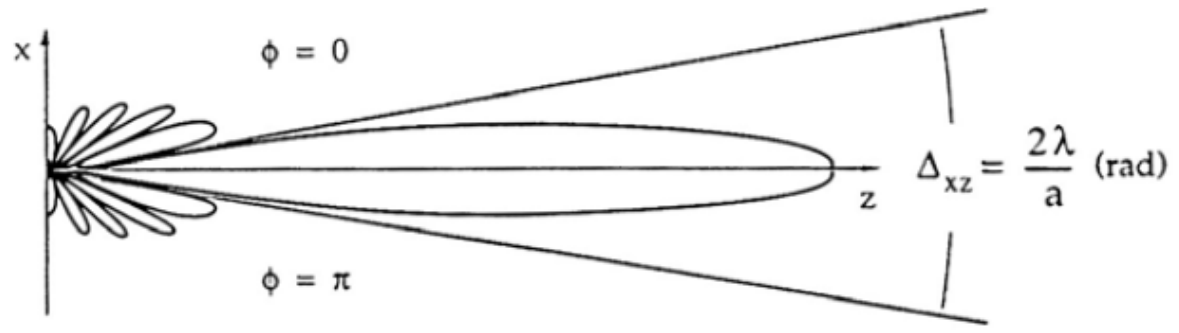
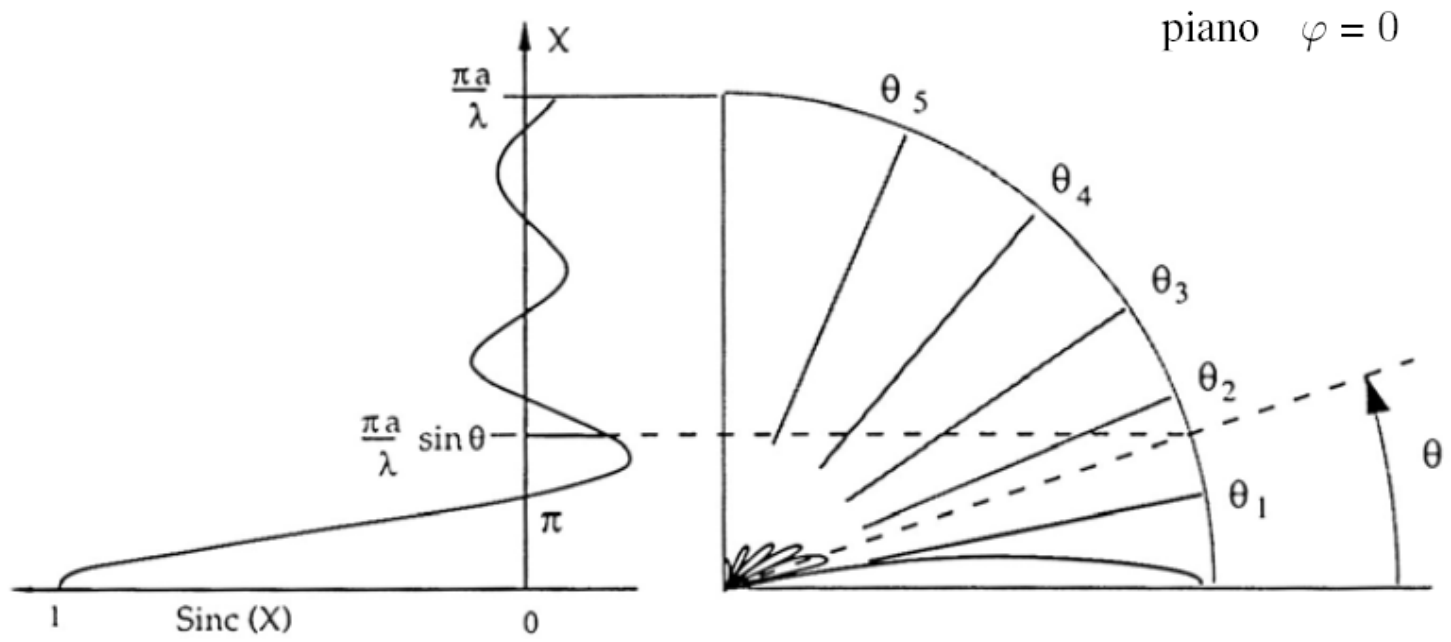
$$\operatorname{sinc}(x) = \frac{\sin x}{x}$$

$$\begin{aligned}
\vec{E} &= j \frac{e^{-jkr}}{2\lambda r} \vec{u}_r \times \vec{L} \\
&= j \frac{e^{-jkr}}{\lambda r} E_o a b \operatorname{sinc}(\pi ua/\lambda) \operatorname{sinc}(\pi vb/\lambda) \sqrt{1 - \sin^2 \vartheta \sin^2 \varphi} \vec{p}
\end{aligned}$$

$$\vec{p} = \frac{\vec{u}_\vartheta \cos \varphi - \vec{u}_\varphi \cos \vartheta \sin \varphi}{\sqrt{1 - \sin^2 \vartheta \sin^2 \varphi}}$$

$$\begin{aligned}
\vec{H} &= -\frac{j}{\eta} \frac{e^{-jkr}}{2\lambda r} \left(\vec{L} - \vec{u}_r (\vec{u}_r \cdot \vec{L}) \right) \\
&= j \frac{e^{-jkr}}{\eta \lambda r} E_o a b \operatorname{sinc}(\pi ua/\lambda) \operatorname{sinc}(\pi vb/\lambda) \sqrt{1 - \sin^2 \vartheta \sin^2 \varphi} \vec{u}_r \times \vec{p}
\end{aligned}$$

$$\begin{aligned}
K &= \frac{1}{8\eta \lambda^2} \left(|\vec{L}|^2 - |\vec{u}_r \cdot \vec{L}|^2 \right) \\
&= \frac{a^2 b^2}{2\eta \lambda^2} |E_o|^2 \operatorname{sinc}^2(\pi ua/\lambda) \operatorname{sinc}^2(\pi vb/\lambda) (1 - \sin^2 \vartheta \sin^2 \varphi)
\end{aligned}$$



Apertura rettangolare illuminata uniformemente con variazione di fase lineare

$$\vec{E}_o = E_o \vec{u}_x e^{-j k(\delta_x x + \delta_y y)}$$

$$\begin{aligned} \vec{L} &= -2 E_o \vec{u}_y \int_{-a/2}^{a/2} e^{j k(u - \delta_x) x'} dx' \int_{-b/2}^{b/2} e^{j k(v - \delta_y) y'} dy' \\ &= -2 E_o a b \operatorname{sinc}(\pi(u - \delta_x)a/\lambda) \operatorname{sinc}(\pi(v - \delta_y)b/\lambda) \vec{u}_y \end{aligned}$$

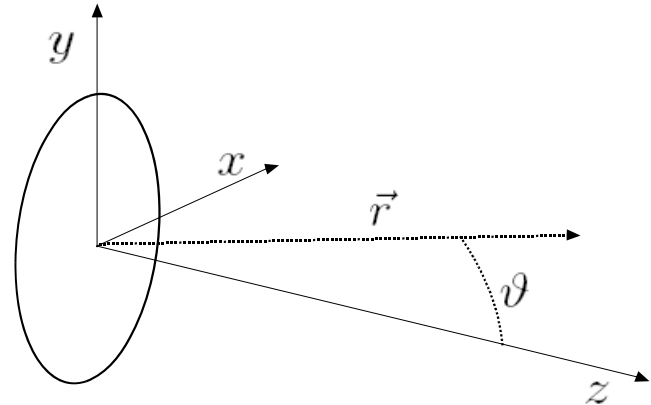
$$\vec{E} = j \frac{e^{-jkr}}{\lambda r} E_o a b \operatorname{sinc}(\pi(u - \delta_x)a/\lambda) \operatorname{sinc}(\pi(v - \delta_y)b/\lambda) \sqrt{1 - \sin^2 \vartheta \sin^2 \varphi} \vec{p}$$

$$\vec{p} = \frac{\vec{u}_\vartheta \cos \varphi - \vec{u}_\varphi \cos \vartheta \sin \varphi}{\sqrt{1 - \sin^2 \vartheta \sin^2 \varphi}}$$

$$\vec{H} = j \frac{e^{-jkr}}{\eta \lambda r} E_o a b \operatorname{sinc}(\pi(u - \delta_x)a/\lambda) \operatorname{sinc}(\pi(v - \delta_y)b/\lambda) \sqrt{1 - \sin^2 \vartheta \sin^2 \varphi} \vec{u}_r \times \vec{p}$$

$$K = \frac{a^2 b^2}{2 \eta \lambda^2} |E_o|^2 \operatorname{sinc}^2(\pi(u - \delta_x)a/\lambda) \operatorname{sinc}^2(\pi(v - \delta_y)b/\lambda) (1 - \sin^2 \vartheta \sin^2 \varphi)$$

Apertura circolare illuminata uniformemente $\vec{E}_o = E_o \vec{u}_x$



$$\begin{aligned}
 \vec{L} &= 2 \int_{\text{apertura}} e^{j k \vec{u}_r \cdot \vec{r}'} M_s(\vec{r}') ds' \\
 &= -2 E_o \vec{u}_y \int_0^{2\pi} \int_0^a e^{j k \rho \sin \vartheta \cos(\varphi - \varphi')} \rho d\rho d\varphi' \\
 &= -4 \pi a^2 E_o \frac{J_1(k a \sin \vartheta)}{k a \sin \vartheta} \vec{u}_y
 \end{aligned}$$

$$\begin{aligned}
\vec{E} &= j \frac{e^{-jkr}}{2\lambda r} \vec{u}_r \times \vec{L} \\
&= j \frac{e^{-jkr}}{\lambda r} E_o 2\pi a^2 \frac{J_1(ka \sin \vartheta)}{ka \sin \vartheta} \sqrt{1 - \sin^2 \vartheta \sin^2 \varphi} \vec{p}
\end{aligned}$$

$$\vec{p} = \frac{\vec{u}_\vartheta \cos \varphi - \vec{u}_\varphi \cos \vartheta \sin \varphi}{\sqrt{1 - \sin^2 \vartheta \sin^2 \varphi}}$$

$$\begin{aligned}
\vec{H} &= -\frac{j}{\eta} \frac{e^{-jkr}}{2\lambda r} \left(\vec{L} - \vec{u}_r (\vec{u}_r \cdot \vec{L}) \right) \\
&= j \frac{e^{-jkr}}{\eta \lambda r} E_o 2\pi a^2 \frac{J_1(ka \sin \vartheta)}{ka \sin \vartheta} \sqrt{1 - \sin^2 \vartheta \sin^2 \varphi} \vec{u}_r \times \vec{p}
\end{aligned}$$

$$\begin{aligned}
K &= \frac{1}{8\eta \lambda^2} \left(|\vec{L}|^2 - |\vec{u}_r \cdot \vec{L}|^2 \right) \\
&= \frac{2\pi^2 a^4}{\eta \lambda^2} |E_o|^2 \left(\frac{J_1(ka \sin \vartheta)}{ka \sin \vartheta} \right)^2 (1 - \sin^2 \vartheta \sin^2 \varphi)
\end{aligned}$$